Weighted approximation by q-Ibragimov-Gadjiev operators

SERAP HERDEM^{1,*}AND İBRAHIM BÜYÜKYAZICI²

 ¹ Department of Mechanical Engineering, Piri Reis University, Postane Mahallesi, Eflatun Sk. No:8, 34 940 Tuzla, Istanbul, Turkey
 ² Department of Mathematics, Faculty of Science, Ankara University, Dögol Caddesi, 06 100, Tandoğan, Ankara, Turkey

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Abstract. The present paper deals with a q-generalization of Ibragimov-Gadjiev operators, which were constructed by Ibragimov and Gadjiev in 1970. We study convergence properties on the interval $[0, \infty)$ and obtain the rate of convergence in terms of a weighted modulus of continuity. We also give some representation formulas of the operators using q-derivatives.

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1. Introduction

In 1970, İbragimov and Gadjiev [10] introduced the following general sequence of positive linear operators defined on the space C[0, A]:

$$L_n\left(f;x\right) = \sum_{\nu=0}^{\infty} f\left(\frac{\nu}{n^2 \Psi_n\left(0\right)}\right) \left[\frac{\partial^{\nu}}{\partial u^{\nu}} K_n\left(x,t,u\right)|_{\substack{u=\alpha_n \Psi_n\left(t\right)\\t=0}}\right] \frac{(-\alpha_n \Psi_n\left(0\right))^{\nu}}{\nu!} \quad (1)$$

which contains the well-known operators of Bernstein, Bernstein-Chlodowsky, Szász and Baskakov as a particular case. Over time, these operators, called Ibragimov-Gadjiev operators, have been deeply investigated and different generalizations have been obtained in numerous papers, see [1 - 4, 8, 9, 11] and the literature cited therein.

In [9], an extension in q-Calculus of Ibragimov-Gadjiev operators was constructed. Here we briefly describe some details. First of all, we should mention some basic definitions of q-Calculus. We refer to [13] as a good guide. Let q > 0. For any natural number $n \in \mathbb{N} \cup \{0\}$, the q-integer $[n] = [n]_q$ is defined by:

$$[n] = \frac{1 - q^n}{1 - q}, \qquad [0] = 0$$

and the q-factorial $[n]! = [n]_q!$ by

$$[n]! = [1] [2] \cdots [n], \qquad [0]! = 1.$$

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^{*}Corresponding author. *Email addresses:* skaya@pirireis.edu.tr (S. Herdem), ibuyukyazici@gmail.com (İ. Büyükyazıcı)

For integers $0 \le k \le n$, the q-binomial coefficient is defined by:

$$\begin{bmatrix} n\\k \end{bmatrix} = \frac{[n]!}{[n-k]!\,[k]!}.$$

Clearly, for q = 1,

$$[n]_1 = n, \quad [n]_1! = n!, \quad \begin{bmatrix} n\\k \end{bmatrix}_1 = \begin{pmatrix} n\\k \end{pmatrix}.$$

The q-derivative of a function $f : \mathbb{R} \to \mathbb{R}$ is defined by:

$$D_q f(x) = \begin{cases} \frac{f(x) - f(qx)}{(1-q)x}, & x \neq 0\\ f'(0), & x = 0 \end{cases},$$

for functions which are differentiable at x = 0 and the higher q-derivatives are

$$D_q^0 f = f, \ D_q^n f = D_q \left(D_q^{n-1} f \right), \ n = 1, 2, 3, \dots$$

The product rule is

$$D_{q}(f(x)g(x)) = D_{q}f(x)g(qx) + D_{q}g(x)f(x).$$
(2)

The q-analogue of $(x-a)^n$ is the polynomial

$$(x-a)_q^n = \begin{cases} 1, & n=0\\ (x-a) (x-qa) \cdots (x-q^{n-1}a), & n \ge 1 \end{cases}$$

The q-exponential functions are given as follows:

$$e_q(x) = \sum_{k=0}^{\infty} \frac{x^k}{[k]!} = \frac{1}{\left(1 - (1 - q)x\right)_q^{\infty}}, \ |x| < \frac{1}{1 - q}, \ |q| < 1,$$
(3)

$$E_q(x) = \sum_{k=0}^{\infty} q^{\frac{k(k-1)}{2}} \frac{x^k}{[k]!} = (1 + (1-q)x)_q^{\infty}, \ x \in \mathbb{R}, \ |q| < 1.$$
(4)

It is clear from equations (3) and (4) that

$$e_q(x) E_q(-x) = 1.$$

The following q-generalization of Taylor's formula was introduced by Jackson (see [6]).

Theorem 1. If the function f can be expanded in a convergent power series and if q is not a root of unity, then

$$f(x) = \sum_{n=0}^{\infty} D_q^n f(a) \, \frac{(x-a)_q^n}{[n]!}.$$

In [9], the authors introduced a q-analogue of the Ibragimov-Gadjiev operators that are called the q-Ibragimov-Gadjiev operators as follows:

Let 0 < q < 1 and let A > 0 be a given number and $\{\psi_n(t)\}$ a sequence of functions in C[0, A] such that $\psi_n(t) > 0$ for each $t \in [0, A]$, $\lim_{n \to \infty} \frac{1}{[n]^2 \psi_n(0)} = 0$. Also, let $\{\alpha_n\}$ be a sequence of positive real numbers such that $\lim_{n \to \infty} \frac{\alpha_n}{[n]} = 1$. Assume that a function $K^q_{-n}(x, t, x)$ ($x, t \in [0, A]$).

Assume that a function $K_{n,\nu}^q(x,t,u)$ $(x,t \in [0,A], -\infty < u < \infty)$ depending on the three parameters n, ν and q satisfies the following conditions:

- (i) The function $K_{n,\nu}^q(x,t,u)$ is infinitely q-differentiable with respect to u for fixed $x, t \in [0, A]$,
- (*ii*) For any $x \in [0, A]$ and for all $n \in \mathbb{N}$,

$$\sum_{\nu=0}^{\infty} q^{\frac{\nu(\nu-1)}{2}} \left[D_{q,u}^{\nu} K_{n,\nu}^{q}\left(x,t,u\right)|_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} \right] \frac{(-q\alpha_{n}\psi_{n}\left(0\right))^{\nu}}{[\nu]!} = 1$$

(*iii*)
$$\left\{ (-1)^{\nu} D_{q,u}^{\nu} K_{n,\nu}^{q}(x,t,u) |_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} \right\} \ge 0 \ (\nu, n=1,2,\ldots,x \in [0,A]),$$

(This notation means that if the q-derivative with respect to u is taken ν times, then one sets $u = \alpha_n \psi_n(t)$ and t = 0.)

$$(iv) \quad D_{q,u}^{\nu} K_{n,\nu}^{q}\left(x,t,u\right)|_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} = -\left[n\right] xq^{1-\nu} \left[D_{q,u}^{\nu-1} K_{n+m,\nu-1}^{q}\left(x,t,u\right)|_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} \right]$$

 $(\nu, n = 1, 2, \dots, x \in [0, A])$, where m is a number such that m + n is zero or a natural number.

According to these conditions, the q-Ibragimov-Gadjiev operators have the following form:

$$L_{n}\left(f;q;x\right) = \sum_{\nu=0}^{\infty} q^{\frac{\nu(\nu-1)}{2}} f\left(\frac{[\nu]}{[n]^{2}\psi_{n}\left(0\right)}\right) \left[D_{q,u}^{\nu}K_{n,\nu}^{q}\left(x,t,u\right)|_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}}\right] \frac{(-q\alpha_{n}\psi_{n}\left(0\right))^{\nu}}{[\nu]!}$$
(5)

for $x \in \mathbb{R}_+$ and any function f defined on \mathbb{R}_+ .

By ν -multiple application of property (iv), $L_n(f;q;x)$ can be reduced to the form:

$$L_{n}(f;q;x) = \sum_{\nu=0}^{\infty} f\left(\frac{[\nu]}{[n]^{2}\psi_{n}(0)}\right) \frac{[n][n+m][n+2m]\cdots[n+(\nu-1)m]}{[\nu]!} \times (xq\alpha_{n}\psi_{n}(0))^{\nu}K_{n+\nu m,0}^{q}(x,0,\alpha_{n}\psi_{n}(0)).$$
(6)

Note that for q = 1, the operators $L_n(f; q; .)$ are classical Ibragimov-Gadjiev operators. In [9], the authors showed that by taking $q_n \in (0, 1)$ such that $\lim_{n \to \infty} q_n = 1$, $L_n(f; q; .)$ satisfies the conditions of the Bohman-Korovkin theorem, and obtained the other convergence properties on the space C[0, A]. Furthermore, they showed that in special cases, this sequence of operators consist of *q*-positive linear operators. For instance, by choosing $K_{n,\nu}^q(x,t,u) = \left(1 - \frac{q^{1-\nu}ux}{1+t}\right)_q^n$ and $\alpha_n = \frac{[n]}{q}$, $\psi_n(0) = \frac{1}{[n]}$, the operators defined by (5) are transformed into q-Bernstein polynomials which were introduced by Phillips [15], for $\alpha_n = \frac{[n]}{q}$, $\psi_n(0) = \frac{1}{[n]b_n}$ $(\lim_{n \to \infty} b_n = \infty, \lim_{n \to \infty} \frac{b_n}{[n]} = 0)$ the operators become q-Chlodowsky polynomials defined by Karsh and Gupta [14]. For $K_{n,\nu}^q(x,t,u) = E_q\left(-\left[n\right]\left(t+q^{1-\nu}ux\right)\right)$, $\alpha_n = \frac{[n]}{q}, \ \psi_n(0) = \frac{1}{[n]b_n} \left(\lim_{n \to \infty} b_n = \infty, \lim_{n \to \infty} \frac{b_n}{[n]} = 0\right), \ \text{the } q\text{-analogue of the classical Szász-Mirakjan operators, which were defined by Aral [5], can be obtained.}$

The aim of this paper is to establish Korovkin type theorems for continuous and unbounded functions defined on $[0,\infty)$ by the q-Ibragimov-Gadjiev operators. We also give some representations of these operators using the q-differences and divided differences.

2. Preliminary results

We define a suitable q-difference operator as follows:

$$\Delta_q^0 f_{\nu} = f_{\nu}, \ \Delta_q^r f_{\nu} = \Delta_q^{r-1} f_{\nu+1} - q^{r-1} \Delta_q^{r-1} f_{\nu} \ , \ r \ge 1,$$
(7)

where $f_{\nu} = f\left(\frac{[\nu]}{[n]^2\psi_n(0)}\right), \nu \in \mathbb{N}_0.$ Now, suppose that in addition to conditions (i) - (iv) function $K_{n,\nu}^q(x,t,u)$ satisfies

(v) $K_{n,\nu}^{q}(x,t,u)$ is infinitely q- differentiable with respect to x for fixed $t\in[0,A]$,

$$D_{q,x}K_{n,\nu}^{q}(x,t,u)|_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} = -[n]\alpha_{n}\psi_{n}(0)q^{1-\nu}K_{n+m,\nu-1}^{q}(x,0,\alpha_{n}\psi_{n}(0)),$$

where m is the natural number defined in (iv),

(vi) $K_{n,\nu}^q(0,0,u) = 1$ for all n, ν and $u \in \mathbb{R}$.

In what follows we take another look at our operators.

Theorem 2. Let $q \in (0,1)$. The generalized g-Ibraqimov-Gadjiev operators may be expressed in the form:

$$L_{n}(f;q;x) = \sum_{\nu=0}^{\infty} \triangle_{q}^{\nu} f_{0} \frac{[n][n+m]\cdots[n+(\nu-1)m]}{[\nu]!} (q\alpha_{n}\psi_{n}(0))^{\nu} \times K_{n+\nu m,-\nu}^{q}(0,0,\alpha_{n}\psi_{n}(0)) x^{\nu},$$

 $x \in [0, A]$, where $\triangle_a^{\nu} f_0$ is defined as in (7).

Proof. Applying a q-derivative operator to (6) and taking into account product rule (2) and property (v), we have:

$$D_{q,x}L_n(f;q;x) = \sum_{\nu=0}^{\infty} \triangle_q^1 f_{\nu} \frac{[n][n+m]\cdots[n+\nu m]}{[\nu]!} (q\alpha_n\psi_n(0))^{\nu+1} K_{n+(\nu+1)m,-1}^q(x,0,\alpha_n\psi_n(0)) x^{\nu}$$

for $n \in \mathbb{N}$ and $x \in [0, A]$. By induction with respect to r we can prove:

$$D_{q,x}^{r}L_{n}(f;q;x) = \sum_{\nu=0}^{\infty} \triangle_{q}^{r} f_{\nu} \frac{[n][n+m]\cdots[n+(\nu+r-1)m]}{[\nu]!} (q\alpha_{n}\psi_{n}(0))^{\nu+r} \times K_{n+(\nu+r)m,-r}^{q}(x,0,\alpha_{n}\psi_{n}(0)) x^{\nu}.$$

From the expansion of the above series by choosing x = 0, we observe that all terms except the first one are zero,

$$D_{q,x}^{r}L_{n}(f;q;0) = \triangle_{q}^{r}f_{0}\frac{[n][n+m]\cdots[n+(r-1)m]}{[0]!}(q\alpha_{n}\psi_{n}(0))^{r} \times K_{n+rm,-r}^{q}(0,0,\alpha_{n}\psi_{n}(0)), r = 0, 1, 2, \dots$$

Hence, choosing a = 0 in Theorem 1, we obtain

$$L_{n}(f;q;x) = \sum_{\nu=0}^{\infty} \triangle_{q}^{\nu} f_{0} \frac{[n][n+m] \cdots [n+(\nu-1)m]}{[\nu]!} (q\alpha_{n}\psi_{n}(0))^{\nu} \times K_{n+(\nu)m,-\nu}^{q}(0,0,\alpha_{n}\psi_{n}(0)) x^{\nu},$$
(8)

which completes the proof.

Now, we will give another representation of operators (5) in terms of divided differences, using the following theorem given in [16, p. 44]. Here, $[x_0, x_1, \ldots, x_k; f]$ denotes the divided difference of the function f with respect to distinct points in the domain of f,

$$[x_0; f] = f(x_0), [x_0, x_1; f] = \frac{f(x_1) - f(x_0)}{x_1 - x_0}, \dots,$$
$$[x_0, x_1, \dots, x_k; f] = \frac{[x_1, \dots, x_k; f] - [x_0, \dots, x_{k-1}; f]}{x_k - x_0}.$$

Theorem 3. For all $j, k \ge 0$,

$$f[x_j, x_{j+1}, \dots, x_{j+k}] = \frac{\Delta_q^k f_j}{q^{k(2j+k-1)/2} [k]!},$$

where $x_j = [j]$ and $[k]! = [k] [k-1] \cdots [1]$.

Corollary 1. The q-Ibragimov-Gadjiev operators can be represented as:

$$L_{n}(f;q;x) = \sum_{\nu=0}^{\infty} \frac{(q\alpha_{n}x)^{\nu}}{[n]^{2\nu}} \left[0, \frac{[1]}{[n]^{2}\psi_{n}(0)}, \frac{[2]}{[n]^{2}\psi_{n}(0)}, \dots, \frac{[\nu]}{[n]^{2}\psi_{n}(0)}; f \right] \\ \times q^{\frac{(\nu-1)\nu}{2}} [n][n+m] \cdots [n+(\nu-1)m] K_{n+\nu m,-\nu}^{q}(0,0,\alpha_{n}\psi_{n}(0)).$$

Proof. The proof is obvious from equality (8) and Theorem 3.

Remark 1. It will be noted that for q = 1, this is the result obtained by Ibragimov and Gadjiev [10] for the operators defined by (1).

Lemma 1. For the operators defined by (5), the following equalities hold:

$$L_n(1;q;x) = 1, \tag{9}$$

$$L_n(t;q;x) = \frac{q\alpha_n}{[n]}x,$$
(10)

$$L_n(t^2;q;x) = \left(\frac{q\alpha_n}{[n]}x\right)^2 \frac{q[n+m]}{[n]} + \left(\frac{q\alpha_n}{[n]}x\right) \frac{1}{[n]^2\psi_n(0)},$$
(11)

$$L_n\left(t^3;q;x\right) = \frac{q^6\alpha_n^3\left[n+m\right]\left[n+2m\right]x^3}{\left[n\right]^5} + \frac{\left(q^4+2q^3\right)\alpha_n^2\left[n+m\right]x^2}{\left[n\right]^5\psi_n\left(0\right)} + \frac{q\alpha_n x}{\left[n\right]^5\psi_n^2\left(0\right)},\tag{12}$$

$$L_{n}\left(t^{4};q;x\right) = \frac{q^{10}\alpha_{n}^{4}\left[n+m\right]\left[n+2m\right]\left[n+3m\right]x^{4}}{\left[n\right]^{7}} + \frac{\left(q^{8}+2q^{7}+3q^{6}\right)\alpha_{n}^{3}\left[n+m\right]\left[n+2m\right]x^{3}}{\left[n\right]^{7}\psi_{n}\left(0\right)} + \frac{\left(q^{5}+3q^{4}+3q^{3}\right)\alpha_{n}^{2}\left[n+m\right]x^{2}}{\left[n\right]^{7}\psi_{n}^{2}\left(0\right)} + \frac{q\alpha_{n}x}{\left[n\right]^{7}\psi_{n}^{3}\left(0\right)}.$$
(13)

Proof. For s = 0, 1, 2, the equalities of $L_n(t^s; q; .)$ were obtained in view of the definition of the operators defined by (5) in [9]. So this part of proof is omitted. Hence for s = 3, 4, we will prove this lemma in terms of q-differences. Using equality (7) for s = 3 one has

$$\begin{split} & \bigtriangleup_q^0 f_0 = f_0 = 0 \\ & \bigtriangleup_q^1 f_0 = f_1 - f_0 = \left(\frac{1}{[n]^2 \psi_n(0)}\right)^3 \\ & \bigtriangleup_q^2 f_0 = f_2 - (1+q) f_1 + q f_0 = \left(\frac{[2]}{[n]^2 \psi_n(0)}\right)^3 - (1+q) \left(\frac{1}{[n]^2 \psi_n(0)}\right)^3 \\ & \bigtriangleup_q^3 f_0 = f_3 - (1+q+q^2) f_2 + (q+q^2+q^3) f_1 - q^3 f_0 \\ & = \left(\frac{[3]}{[n]^2 \psi_n(0)}\right)^3 - (1+q+q^2) \left(\frac{[2]}{[n]^2 \psi_n(0)}\right)^3 \\ & + (q+q^2+q^3) \left(\frac{1}{[n]^2 \psi_n(0)}\right)^3. \end{split}$$

By using Theorem 2, we can write:

$$\begin{split} L_n\left(t^3;q;x\right) = & \bigtriangleup_q^0 f_0 + \bigtriangleup_q^1 f_0 \frac{[n]}{[1]!} q \alpha_n \psi_n\left(0\right) K_{n+m,-1}^q\left(0,0,\alpha_n \psi_n\left(0\right)\right) x + \\ & + \bigtriangleup_q^2 f_0 \frac{[n][n+m]}{[2]!} (q \alpha_n \psi_n\left(0\right))^2 K_{n+2m,-2}^q\left(0,0,\alpha_n \psi_n\left(0\right)\right) x^2 + \\ & + \bigtriangleup_q^3 f_0 \frac{[n][n+m][n+2m]}{[3]!} (q \alpha_n \psi_n\left(0\right))^3 K_{n+3m,-3}^q\left(0,0,\alpha_n \psi_n\left(0\right)\right) x^3. \end{split}$$

Taking into account condition (vi), by a straightforward calculation, we obtain equality (12).

We now prove (13). For s = 4, we have:

$$\begin{split} & \triangle_q^0 f_0 = f_0 = 0 \\ & \triangle_q^1 f_0 = f_1 - f_0 = \left(\frac{1}{[n]^2 \psi_n(0)}\right)^4 \\ & \triangle_q^2 f_0 = f_2 - (1+q) f_1 + q f_0 = \left(\frac{[2]}{[n]^2 \psi_n(0)}\right)^4 - (1+q) \left(\frac{1}{[n]^2 \psi_n(0)}\right)^4 \\ & \triangle_q^3 f_0 = f_3 - (1+q+q^2) f_2 + (q+q^2+q^3) f_1 - q^3 f_0 \\ & = \left(\frac{[3]}{[n]^2 \psi_n(0)}\right)^4 - (1+q+q^2) \left(\frac{[2]}{[n]^2 \psi_n(0)}\right)^4 + \\ & + (q+q^2+q^3) \left(\frac{1}{[n]^2 \psi_n(0)}\right)^4 \\ & \Delta_q^4 f_0 = f_4 - (1+q+q^2+q^3) f_3 + (q+q^2+2q^3+q^4+q^5) f_2 \\ & - (q^3+q^4+q^5+q^6) f_1 + q^6 f_0 \\ & = \left(\frac{[4]}{[n]^2 \psi_n(0)}\right)^4 - (1+q+q^2+q^3) \left(\frac{[3]}{[n]^2 \psi_n(0)}\right)^4 \\ & + (q+q^2+2q^3+q^4+q^5) \left(\frac{[2]}{[n]^2 \psi_n(0)}\right)^4 - (q^3+q^4+q^5+q^6) \left(\frac{1}{[n]^2 \psi_n(0)}\right) \end{split}$$

Thus applying again Theorem 2 and using condition (vi), proceeding similarly, we get (13).

3. Weighted approximation

This section is devoted to obtaining weighted approximation properties of the q-Ibragimov-Gadjiev operators. Note that we will use a weighted Korovkin type theorem proved by Gadjiev [7]. When considering that theorem we should mention some notations: Let $\rho(x) = 1 + x^2$ be a weight function and $B_{\rho}[0,\infty)$ the set of all functions f defined on the semi-axis $[0,\infty)$ satisfying the condition $|f(x)| \leq M_f \rho(x)$, where M_f is a constant depending only on f. $C_{\rho}[0,\infty)$ denotes the subspace of all continuous functions belonging to $B_{\rho}[0,\infty)$ and $C_{\rho}^k[0,\infty)$ denotes the subspace of

all functions $f \in C_{\rho}[0,\infty)$ with $\lim_{|x|\to\infty} \frac{f(x)}{\rho(x)} = k_f < \infty$. Obviously $C_{\rho}^k[0,\infty)$ is a linear normed space with the ρ -norm:

$$||f||_{\rho} = \sup_{x \in [0,b_n]} \frac{|f(x)|}{\rho(x)},$$

where $\{b_n\}$ is a sequence of positive numbers, which has a finite or an infinite limit.

Theorem 4 (see [7]). Let $\{T_n\}$ be a sequence of linear positive operators which are mappings from C_{ρ} into B_{ρ} satisfying the conditions:

$$\lim_{n \to \infty} \left\| T_n(t^k; x) - x^k \right\|_{\rho} = 0, \ k = 0, 1, 2.$$

Then, for any function $f \in C^k_{\rho}$,

$$\lim_{n \to \infty} \|T_n f - f\|_{\rho} = 0.$$

In order to study convergence properties of operators (5) into the weighted approximation fields we consider the following modified form of the *q*-Ibragimov-Gadjiev operators.

Let $\{b_n\}$ be a sequence of positive numbers, which has a finite or an infinite limit and $\{\psi_n(t)\}\$ be a sequence of functions in $C[0, b_n]$ such that $\psi_n(t) > 0$ for each $t \in [0, b_n]$, satisfying the condition:

$$\lim_{n \to \infty} \frac{1}{[n]^2 \psi_n(0)} = 0.$$
(14)

Also, let $\{\alpha_n\}$ be a sequence of positive real numbers such that

$$\frac{\alpha_n}{[n]} = 1 + O\left(\frac{1}{\left[n\right]^2 \psi_n\left(0\right)}\right). \tag{15}$$

Now, assume that a function $K_{n,\nu}^q(x,t,u)$ $(x,t \in [0,b_n], -\infty < u < \infty)$ depending on the three parameters n, ν and q satisfies the following conditions:

- (*i^o*) The function $K_{n,\nu}^q(x,t,u)$ is infinitely *q*-differentiable with respect to *u* for fixed $x, t \in [0, b_n]$,
- (ii^{o}) For any $x \in [0, b_n]$ and for all $n \in \mathbb{N}$,

$$\sum_{\nu=0}^{\infty} q^{\frac{\nu(\nu-1)}{2}} \left[D_{q,u}^{\nu} K_{n,\nu}^{q} \left(x, t, u \right) |_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} \right] \frac{(-q\alpha_{n}\psi_{n}(0))^{\nu}}{[\nu]!} = 1,$$

$$(iii^{o}) \left\{ (-1)^{\nu} D_{q,u}^{\nu} K_{n,\nu}^{q} (x,t,u) |_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} \right\} \ge 0 (\nu, n = 1, 2, \dots, x \in [0, b_{n}]),$$

 $(iv^{o}) \quad D_{q,u}^{\nu} K_{n,\nu}^{q} (x,t,u) \mid_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} = - [n] xq^{1-\nu} \left[D_{q,u}^{\nu-1} K_{n+m,\nu-1}^{q} (x,t,u) \mid_{\substack{u=\alpha_{n}\psi_{n}(t)\\t=0}} \right]$ $(\nu, n = 1, 2, \dots, x \in [0, b_{n}]), \text{ where } m \text{ is a number such that } m+n \text{ is zero or a}$

 $(\nu, n = 1, 2, ..., x \in [0, b_n])$, where m is a number such that m + n is zero or natural number.

Under above conditions we call operators (5) $L_n^*(f;q;.)$. It is clear that $L_n^*(f;q;.)$ holds for the equalities in Lemma 1. In the case $\lim_{n\to\infty} b_n = A$ and $\lim_{n\to\infty} \frac{\alpha_n}{[n]} = 1$ we obtain q-İbragimov-Gadjiev operators defined by (5).

Theorem 5. Let $q = q_n$, $0 < q_n < 1$, satisfy the condition $(q_n) \to 1$ as $n \to \infty$. Then for each function $f \in C_{\rho}^k[0,\infty)$,

$$\lim_{n \to \infty} \|L_n^*(f;q;.) - f\|_{\rho,[0,b_n]} = 0.$$

Proof. Making use of the Korovkin type theorem on weighted approximation, it is sufficient to verify the following three conditions:

$$\lim_{n \to \infty} \left\| L_n^*(t^k; q; x) - x^k \right\|_{\rho} = 0, \ k = 0, 1, 2.$$
(16)

From equality (9) clearly $||L_n(1,q_n,x)-1||_{\rho} \to 0$ as $n \to \infty$ on $[0,b_n]$. From equalities (10) and (14) we have:

$$\sup_{x \in [0,b_n)} \frac{\left|L_n^*\left(t;q;x\right) - x\right|}{1 + x^2} \le \left|\frac{q\alpha_n}{[n]} - 1\right| \sup_{x \in [0,b_n)} \frac{x}{1 + x^2}$$
$$\le \left|\frac{q\alpha_n}{[n]} - 1\right|,$$

hence the condition in (16) holds for k = 1.

Similarly, using equality (11) for k = 2 we can write:

$$\sup_{x \in [0,b_n)} \frac{\left|L_n^*\left(t^2; q; x\right) - x^2\right|}{1 + x^2} \le \left| \left(\frac{q\alpha_n}{[n]}\right)^2 \frac{q\left[n+m\right]}{[n]} - 1 \right| \sup_{x \in [0,b_n)} \frac{x^2}{1 + x^2} \\ + \frac{q\alpha_n}{[n]} \frac{1}{[n]^2 \psi_n\left(0\right)} \sup_{x \in [0,b_n)} \frac{x}{1 + x^2} \\ \le \left| \left(\frac{q\alpha_n}{[n]}\right)^2 \frac{q\left[n+m\right]}{[n]} - 1 \right| + \frac{q\alpha_n}{[n]} \frac{1}{[n]^2 \psi_n\left(0\right)},$$

which implies that (16) also holds for k = 2. Therefore, the desired result follows from Theorem 4.

Now, we will give an estimate concerning the rate of convergence for the q-Ibragimov-Gadjiev operators. As known, the first modulus of continuity $\omega(f; \delta)$ does not tend to zero as $\delta \to 0$ on the infinite interval. For this reason, we consider

$$\Omega(f; \delta) = \sup_{|h| \le \delta, x \ge 0} \frac{|f(x+h) - f(x)|}{(1+h^2)(1+x^2)}$$

a weighted modulus of smoothness associated to the space $C_{\rho}^{k}[0,\infty)$.

 $\Omega(f; \delta)$ possesses the following properties (see [12]):

- $\Omega(f; \delta)$ is a monotonically increasing function of $\delta, \delta \ge 0$,

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- for every $f\in C^{k}_{\rho}\left[0,\infty\right),$ $\lim_{\delta\rightarrow0}\Omega\left(f;\delta
 ight)=0,$
- for each positive value of λ

$$\Omega\left(f;\lambda\delta\right) \le 2\left(1+\lambda\right)\left(1+\delta^{2}\right)\Omega\left(f;\delta\right).$$
(17)

From inequality (17) and the definition of $\Omega(f; \delta)$ we get

$$|f(t) - f(x)| \le 2\left(1 + \frac{|t - x|}{\delta}\right) (1 + \delta^2) \Omega(f; \delta) (1 + x^2) (1 + (t - x)^2), \quad (18)$$

for every $f \in C^k_{\rho}[0,\infty)$ and $x,t \in [0,\infty)$.

Theorem 6. Let $q = q_n$, $0 < q_n < 1$ satisfies the condition $q_n \to 1$ as $n \to \infty$. If $f \in C^k_{\rho}[0,\infty)$, then the inequality

$$\sup_{x \ge 0} \frac{|L_n^*(f;q;x) - f(x)|}{(1+x^2)^3} \le M\Omega\left(f; \left([n]^2 \psi_n(0)\right)^{-1/2}\right)$$

is satisfied for a sufficiently large n, where M is a constant independent of α_n , $\psi_n(0)$.

Proof. From (18) we can write:

$$\begin{split} |L_{n}^{*}\left(f;q;x\right) - f\left(x\right)| &\leq 2\left(1 + \delta_{n}^{2}\right)\Omega\left(f;\delta_{n}\right)\left(1 + x^{2}\right)\sum_{\nu=0}^{\infty}P_{\nu,n}^{q}\left(x\right) \\ &\times \left(1 + \frac{\left|\frac{[\nu]}{[n]^{2}\psi_{n}(0)} - x\right|}{\delta_{n}}\right)\left(1 + \left(\frac{[\nu]}{[n]^{2}\psi_{n}\left(0\right)} - x\right)^{2}\right) \\ &\leq 4\left(1 + x^{2}\right)\Omega\left(f;\delta_{n}\right)\left\{1 + \frac{1}{\delta_{n}}\sum_{\nu=0}^{\infty}\left|\frac{[\nu]}{[n]^{2}\psi_{n}\left(0\right)} - x\right|P_{\nu,n}^{q}\left(x\right) \\ &+ \sum_{\nu=0}^{\infty}\left(\frac{[\nu]}{[n]^{2}\psi_{n}\left(0\right)} - x\right)^{2}P_{\nu,n}^{q}\left(x\right) \\ &+ \frac{1}{\delta_{n}}\sum_{\nu=0}^{\infty}\left|\frac{[\nu]}{[n]^{2}\psi_{n}\left(0\right)} - x\right|\left(\frac{[\nu]}{[n]^{2}\psi_{n}\left(0\right)} - x\right)^{2}P_{\nu,n}^{q}\left(x\right)\right\}, \end{split}$$

where $P_{\nu,n}^q(x) = q^{\frac{\nu(\nu-1)}{2}} \left[D_{q,u}^{\nu} K_{n,\nu}^q(x,t,u) |_{\substack{u=\alpha_n\psi_n(t)\\t=0}} \right] \frac{(-q\alpha_n\psi_n(0))^{\nu}}{[\nu]!}$ and $\delta_n > 0$. Applying the Cauchy-Schwarz inequality we obtain

$$|L_n^*(f;q;x) - f(x)| \le 4\left(1 + x^2\right)\Omega(f;\delta_n)\left(1 + \frac{2}{\delta_n}\sqrt{K_1} + K_1 + \frac{1}{\delta_n}K_2\right), \quad (19)$$

where

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$$K_{1} = \sum_{\nu=0}^{\infty} \left(\frac{[\nu]}{[n]^{2}\psi_{n}(0)} - x \right)^{2} P_{\nu,n}^{q}(x)$$
$$K_{2} = \sum_{\nu=0}^{\infty} \left(\frac{[\nu]}{[n]^{2}\psi_{n}(0)} - x \right)^{4} P_{\nu,n}^{q}(x).$$

Using Lemma 1, we have

$$K_1 = L_n^* \left((t-x)^2; q; x \right)$$
$$= \left(1 - 2q \left(\frac{\alpha_n}{[n]} \right) + \left(\frac{\alpha_n}{[n]} \right)^2 \frac{q^3 [n+m]}{[n]} \right) x^2 + q \left(\frac{\alpha_n}{[n]} \right) \frac{1}{[n]^2 \psi_n(0)} x$$

and

$$\begin{split} K_2 =& L_n^* \left((t-x)^4 \, ; q ; x \right) \\ &= \left(\left(\frac{\alpha_n}{[n]} \right)^4 \frac{q^{10} \left[n+m \right] \left[n+2m \right] \left[n+3m \right]}{[n]^3} - 4 \left(\frac{\alpha_n}{[n]} \right)^3 \frac{q^6 \left[n+m \right] \left[n+2m \right]}{[n]^2} \right. \\ &+ 6 \left(\frac{\alpha_n}{[n]} \right)^2 \frac{q^3 \left[n+m \right]}{[n]} - 4q \left(\frac{\alpha_n}{[n]} \right) + 1 \right) x^4 \\ &+ \left(\left(\frac{\alpha_n}{[n]} \right)^3 \frac{(q^8 + 2q^7 + 3q^6) \left[n+m \right] \left[n+2m \right]}{[n]^2} \frac{1}{[n]^2 \psi_n \left(0 \right)} \\ &- 4 \left(\frac{\alpha_n}{[n]} \right)^2 \frac{(q^4 + 2q^3) \left[n+m \right]}{[n]} \frac{1}{[n]^2 \psi_n \left(0 \right)} + 6q \left(\frac{\alpha_n}{[n]} \right) \frac{1}{[n]^2 \psi_n \left(0 \right)} \right) x^3 \\ &+ \left(\left(\frac{\alpha_n}{[n]} \right)^2 \frac{(q^5 + 3q^4 + 3q^3) \left[n+m \right]}{[n]} \frac{1}{([n]^2 \psi_n \left(0 \right))^2} - 4q \left(\frac{\alpha_n}{[n]} \right) \frac{1}{([n]^2 \psi_n \left(0 \right))^2} \right) x^2 \\ &+ q \left(\frac{\alpha_n}{[n]} \right) \frac{1}{([n]^2 \psi_n \left(0 \right))^3} x. \end{split}$$

Using condition (15) we can write

$$K_1 = O\left(\frac{1}{[n]^2\psi_n(0)}\right)(x^2 + x)$$

and

$$K_2 = O\left(\frac{1}{[n]^2\psi_n(0)}\right)\left(x^4 + x^3 + x^2 + x\right),\,$$

thus by substituting these equalities (19) becomes

$$\begin{split} |L_n^*\left(f;q;x\right) - f\left(x\right)| &\leq 4\left(1+x^2\right)\Omega\left(f;\delta_n\right) \left\{ 1 + \frac{2}{\delta_n}\sqrt{O\left(\frac{1}{[n]^2\psi_n\left(0\right)}\right)\left(x^2+x\right)} + \\ &+ O\left(\frac{1}{[n]^2\psi_n\left(0\right)}\right)\left(x^2+x\right) + \frac{1}{\delta_n}O\left(\frac{1}{[n]^2\psi_n\left(0\right)}\right)\left(x^4+x^3+x^2+x\right)\right\}. \end{split}$$

By choosing $\delta_n = \left(\left[n \right]^2 \psi_n(0) \right)^{-1/2}$, for sufficiently large *n*'s, we obtain the desired result.

Remark 2. It will be noted that, for q = 1, this is the result obtained by Gadjiev and Ispir [8] for Ibragimov-Gadjiev operators defined by (1).

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